MAGNETIC PENDULUM IN THREE MAGNETIC POTENTIALS

Introduction

This teaching aid is used to demonstrate the concept of chaos in a nonlinear system. You can learn about chaos by taking the course Nonlinear dynamics (YFX1560).

Equations of motion

The top view projection of the pendulum position (x, y) is described by the following nonlinear equations of motion:

$$\begin{cases} \frac{d^2x}{dt^2} + R\frac{dx}{dt} - \sum_{i=1}^3 \frac{x_i - x}{\left(\sqrt{(x_i - x)^2 + (y_i - y)^2 + D^2}\right)^3} + Cx = 0, \\ \frac{d^2y}{dt^2} + R\frac{dy}{dt} - \sum_{i=1}^3 \frac{y_i - y}{\left(\sqrt{(x_i - x)^2 + (y_i - y)^2 + D^2}\right)^3} + Cy = 0, \end{cases}$$
(1)

where R is proportional to the air resistance and overall attenuation, C is proportional to the effects of gravity, the *i*-th magnet is positioned at (x_i, y_i) where i = 1, 2, 3, and D is the distance between the pendulum at rest position (0, 0) and the plane of magnets.

Additionally, we assume that the pendulum length is long compared to the spacing between the magnets. Thus, we may assume for simplicity that the pendulum moves about on the xy-plane rather than on a sphere with a large radius.



Figure 1: Basins of attraction of the magnets shown with the red, blue and yellow colours.

The basin of attraction of system (1) shown in Fig. 1 is a fractal with an infinite complexity. The prediction of the end state of the pendulum for most initial conditions is impossible in practice especially if the integration time exceeds a several multiples of a time period called the Lyapunov time.

Further information for interested students

If students have any further questions or show a genuine interest in general topics of chaotic dynamics, you may direct them to the Author (e-mail: dmitri.kartofelev@taltech.ee, or URL: https://www.tud.ttu.ee/web/dmitri.kartofelev/).