Lecture №11: Feigenbaum's analysis of period doubling, renormalisation, universal limiting function, discrete-time dynamics analysis, the Poincaré section, the Poincaré map, the Lorenz section, attractor reconstruction

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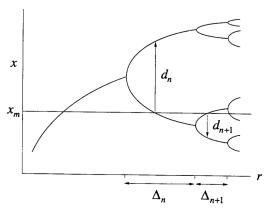
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#### Lecture outline

- Feigenbaum's analysis of period doubling
- The universal route to chaos
- Universal aspects of period doubling in unimodal maps
- Superstable fixed points and period-p orbits
- Renormalisation
- Universal limiting functions and the onset of chaos
- Discrete-time dynamics analysis methods
  - The Poincaré section
  - Return map or the Poincaré map
  - The Lorenz section
- Non-homogenous systems
- Examples studied:
  - The periodically forced damped Duffing oscillator
  - The Rössler attractor

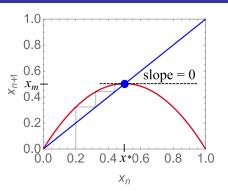
### 1-D unimodal maps and the Feigenbaum constants



$$\delta = \lim_{n \to \infty} \frac{\Delta_{n-1}}{\Delta_n} = \lim_{n \to \infty} \frac{r_{n-1} - r_{n-2}}{r_n - r_{n-1}} \approx 4.669201609... \tag{1}$$

$$\alpha = \lim_{n \to \infty} \frac{d_{n-1}}{d_n} \approx -2.502907875...$$
 (2)

## Superstable fixed point (the logistic map)



$$f(x^*, r) = x^*, \quad x^* = x_m, \quad \text{and} \quad f'(x^*, r) = 0 \quad \Rightarrow r = 2.$$
 (3)

Convergence of  $x_n$  about the non-trivial fixed point  $x^*$ 

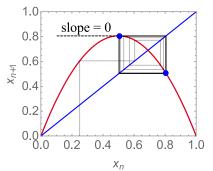
$$\eta_{n+1} = \frac{|f''(x^*, r)|}{2!} \eta_n^2 + O(\eta_n^3), \tag{4}$$

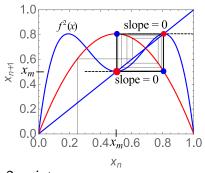
is quadratic. The iterates  $x_n$  converge quadratically.

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## Superstable period-p point (the logistic map)

Superstable period-p orbit:  $x^* = x_m = \max f$  is also a local min. or max. of  $f^p$  map (p-th iterate of map f).



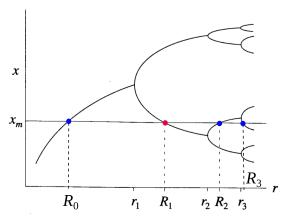


For example in the case of period-2 point

$$f^{2}(x^{*}, r) = x^{*}, \quad (x^{*} = x_{m}) \quad \Rightarrow \quad r = 1 + \sqrt{5}$$
 (5)

$$(f^{2}(x^{*},r))' = \frac{\mathrm{d}}{\mathrm{d}x^{*}}[f(f(x^{*},r))] = f'(f(x^{*},r)) \cdot f'(x^{*},r) = 0.$$
 (6)

### Orbit diagram, superstable period- $2^n$ points



 $r_n$  – stable period- $2^n$  orbit is born (bifurcation point).  $R_n$  – superstable period- $2^n$  point.

$$\lim_{n \to \infty} \frac{R_{n-1}}{R_n} = \delta \tag{7}$$

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### Period doubling bifurcation points

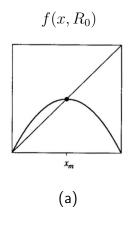
Values of bifurcation points and superstable points of a few first period doublings in the logistic map.

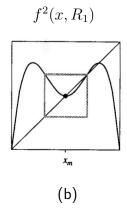
| $r_0=1.0$ non-trivial                | $R_0 = 2.0$                                  | f.p.                 |
|--------------------------------------|--|----------------------|
| $r_1 = 3.0$                          | $R_1 = 1 + \sqrt{5} \approx 3.23607$         | period-2             |
| $r_2 = 1 + \sqrt{6} \approx 3.44949$ | $R_2 \approx 3.49856$                        | period-4             |
| $r_3 \approx 3.54409$                | $R_3 \approx 3.55464$                        | period-8             |
| $r_4 \approx 3.56441$                | $R_4 \approx 3.56667$                        | period-16            |
| $r_5 \approx 3.56875$                | $R_5 \approx 3.56924$                        | period-32            |
| :                                    | :  | :                    |
| $r_{\infty} \approx 3.569945672$     | $R_{\infty} = r_{\infty} \approx 3.56994567$ | period- $2^{\infty}$ |

 $r_{\infty}$  – Onset of chaos (accumulation point).

$$\lim_{n \to \infty} \frac{r_{n-1} - r_{n-2}}{r_n - r_{n-1}} = \lim_{n \to \infty} \frac{R_{n-1}}{R_n} = \delta$$
 (8)

## Feigenbaum's analysis, renormalisation

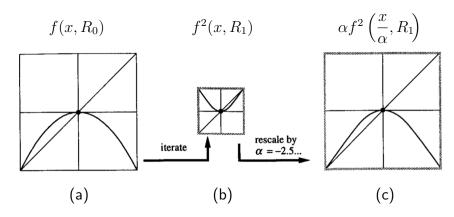






(c)

## Feigenbaum's analysis, renormalisation



**Read:** Mitchell J. Feigenbaum, "The universal metric properties of nonlinear transformations," Journal of Statistical Physics 21(6), pp. 669–706, 1979. Also relevant to the following three slides  $\Downarrow$ .

### Limiting function g(x) and Feigenbaum constant $\alpha$

Let's consider a functional equation in the following form:

$$g(x) = \alpha g(g\left(\frac{x}{\alpha}\right)) = \alpha g^2\left(\frac{x}{\alpha}\right),$$
 (9)

where  $\alpha$  acts as a scaling coefficient and

$$\alpha = \frac{1}{g(1)}. (10)$$

The power series solution is obtained by assuming a power expansion in the following form:

$$g(x) = 1 + ax^{2} + bx^{4} + cx^{6} + dx^{8} + ex^{10} + \dots,$$
 (11)

where the map maximum is quadratic.

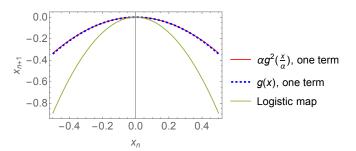
### Limiting function g(x) and Feigenbaum constant $\alpha$

Numeric solution<sup>1</sup>: The one term approximation where

$$g(x) = 1 + ax^2 + O(x^4), (12)$$

results in

$$a = -\frac{1}{2}(1+\sqrt{3}), \quad \alpha = \frac{1}{g(1)} \approx -2.73205, \quad (9.2\% \text{ error}). \quad (13)$$



<sup>&</sup>lt;sup>1</sup>See Mathematica .nb file uploaded to the course webpage.

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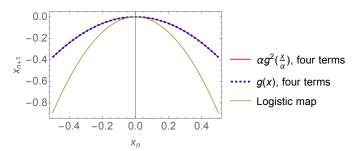
### Limiting function g(x) and Feigenbaum constant $\alpha$

**Numeric solution**<sup>2</sup>: The four term approximation where

$$g(x) = 1 + ax^{2} + bx^{4} + cx^{6} + dx^{8} + O(x^{10}).$$
 (14)

Coefficients:  $a\approx -1.528$ ,  $b\approx 0.1053$ ,  $c\approx 0.02631$ ,  $d\approx -0.003344$  and

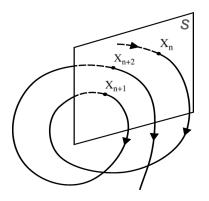
$$\alpha = \frac{1}{g(1)} \approx -2.50316, \quad (0.01\% \text{ error}).$$
 (15)



<sup>&</sup>lt;sup>2</sup>See Mathematica .nb file uploaded to the course webpage.

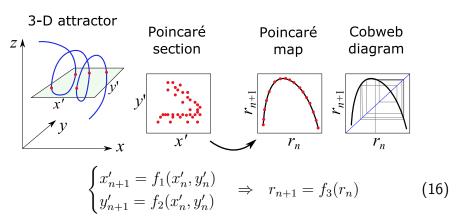
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#### The Poincaré section



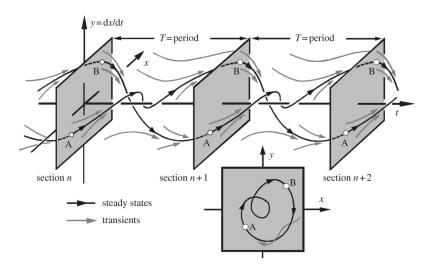
Intersection of an attractor trajectory with hypersurface  ${\cal S}.$ 

#### Discrete-time dynamics analysis



Construction of the Poincaré map  $\vec{P}(\vec{x}') = (f_1(x', y'), f_2(x', y'))^T$  (16). Mapping of the Poincaré section points where r is the radial distance from the origin ("flat" attractor).

# The Poincaré section, periodic forcing<sup>3</sup>



<sup>&</sup>lt;sup>3</sup>Credit: J. Thompson, H. Stewart, 1986, *Nonlinear dynamics and chaos: geometrical methods for engineers and scientists*. Chichester, UK: Wiley.

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#### The Poincaré section, periodic forcing

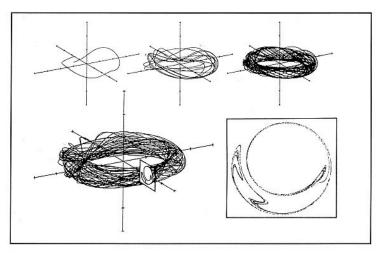


Figure 3. Illustration of a Strange Attractor and the Associated Poincaré Map (from James Gleick, Chaos: Making a New Science [New York: Pengin Books 1987], 143)

### The periodically driven damped Duffing oscillator

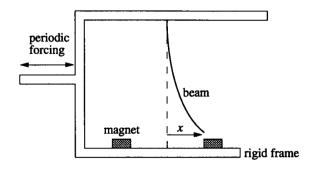


Figure: The mechanical Duffing oscillator.

The non-autonomous equation of motion has the following form:

$$\ddot{x} - x + x^3 + \delta \dot{x} = F \cos \omega t, \tag{17}$$

where  $\delta$  is the damping coefficient, F and  $\omega$  are the forcing strength and frequency, respectively.

## The periodically driven damped Duffing oscillator

We introduce the variable exchange  $y = \dot{x}$  and rewrite (17)

$$\begin{cases} \dot{x} = y, \\ \dot{y} = x - x^3 - \delta y + F \cos \omega t. \end{cases}$$
 (18)

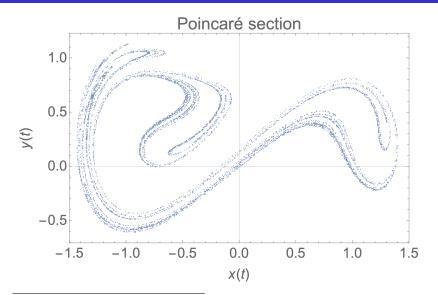
The equivalent 3-D system is obtained by applying variable exchange  $z=\omega t$ 

$$\begin{cases} \dot{x} = y, \\ \dot{y} = x - x^3 - \delta y + F \cos z, \\ \dot{z} = \omega. \end{cases}$$
 (19)

The above holds since

$$z = \int \dot{z} \, dt = \int \omega \, dt = \omega \int dt = \omega t + C.$$
 (20)

### The Poincaré section<sup>4</sup> of the Duffing oscillator



<sup>&</sup>lt;sup>4</sup>See Mathematica .nb file uploaded to the course webpage.

### The Poincaré section dynamics: Duffing oscillator



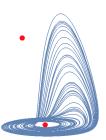
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#### The Rössler attractor

As mentioned in Lecture 9 the Rössler attractor<sup>5</sup> is given by:

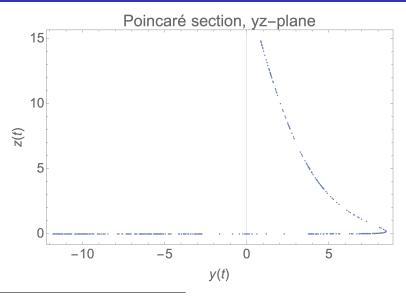
$$\begin{cases} \dot{x} = -y - x, \\ \dot{y} = x + ay, \\ \dot{z} = b + z(x - c). \end{cases}$$
(21)

Chaotic solution exists for a = 0.1, b = 0.1, c = 14.



<sup>&</sup>lt;sup>5</sup>See Mathematica .nb file uploaded to the course webpage.

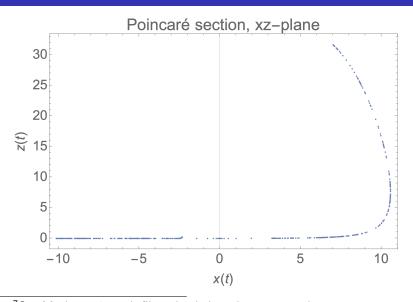
#### The Rössler attractor: the Poincaré sections<sup>6</sup>



<sup>&</sup>lt;sup>6</sup>See Mathematica .nb file uploaded to the course webpage.

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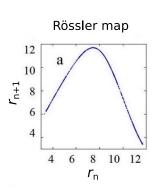
#### The Rössler attractor: the Poincaré sections<sup>7</sup>



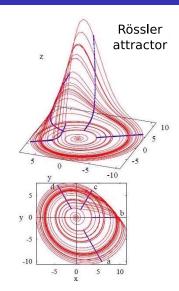
<sup>&</sup>lt;sup>7</sup>See Mathematica .nb file uploaded to the course webpage.

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#### Rössler system: return map and Poincaré sections

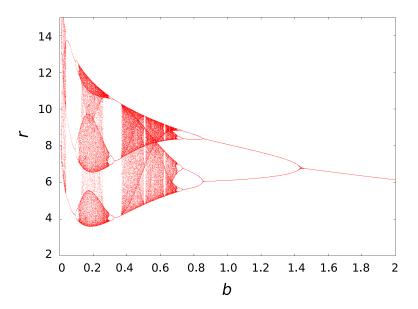


r is the radial distance from the origin



Credit: Y. Maistrenko and R. Paškauskas

#### The Rössler attractor: an orbit diagram



## The Rössler attractor: period doubling



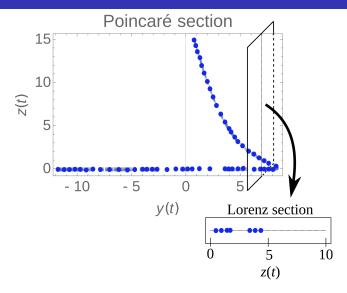
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#### The Rössler attractor: the Poincaré section



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#### The Lorenz section



The Lorenz section — a section of a section.

#### **Conclusions**

- Feigenbaum's analysis of period doubling
- The universal route to chaos
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- Discrete-time dynamics analysis methods
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#### Revision questions

- What are the values of the Feigenbaum constants?
- What are the Feigenbaum constants?
- Define superstable fixed point of a map.
- Define superstable period-p point (or period-p orbit) of a map.
- What are the universals of unimodal maps?
- What is the universal route to chaos?
- Idea behind renormalisation?
- What are the universal limiting functions in the context of maps?
- Name discrete-time dynamics analysis methods.
- What is the Poincaré section?
- What is the Poincaré map (return map)?
- What is the Lorenz section?